

Structured Homotopy Theory from String Theory

Urs Schreiber

(Czech Academy of Science, Prague & NYU, Abu Dhabi)

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Based on joint work with:

V. Braunack-Mayer, D. Fiorenza, J. Huerta, H. Sati,

Abstract.

Homotopy theory is extremely rich, with various structured variants such as equivariant, graded, parameterized and stable homotopy theory. Powerful tools from differential-graded algebra, particularly in the rational approximation, serve for concrete computations.

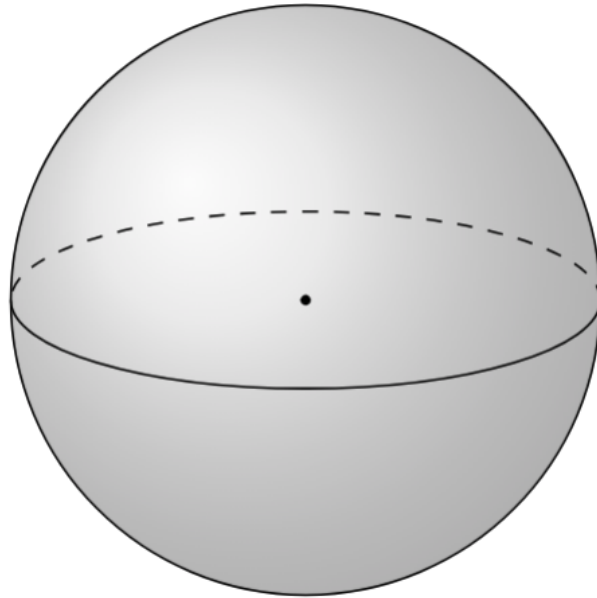
Also string/M-theory is extremely rich, revealing a system of higher dimensional objects (branes) with subtle inter-relations.

In these talks I survey recent insights into a close relation between the two, which provides structured homotopy theory with curious new examples and sheds light on the elusive foundations of string/M-theory.

The 4-sphere will play a surprisingly central role, as initially realized by Hisham Sati.

The 4-sphere

S^4



We will discuss its incarnation
in various flavours of *homotopy theory*:

1)
Super homotopy theory
and
fundamental M2/M5-branes

2)
Parameterized stable homotopy theory
and
gauge enhancement

3)
Equivariant homotopy theory
and
black M-branes

1)
Super homotopy theory
and
fundamental M2/M5-branes

Classical homotopy theory

$$\mathrm{Ho}(\mathrm{Spaces}) := \mathrm{Spaces} \left[\{\text{Isos on all homotopy groups}\}^{-1} \right]$$

see ncatlab.org/nlab/show/Introduction+to+Homotopy+Theory

For example homotopy groups of spheres:

$$\pi_k(S^n) := \mathrm{Hom}_{\mathrm{Ho}(\mathrm{Spaces})} \left(S^k, S^n \right)$$

These exhibit an endlessly rich pattern.

But most of them are torsion groups.

For example for S^4 two of them non-torsion:

generated by the *quaternionic Hopf fibration*:

Quaternionic Hopf fibration

$$\begin{array}{ccc}
 S^7 & \simeq & S(\mathbb{H} \oplus \mathbb{H}) \\
 \downarrow \text{Hopf}_{\mathbb{H}} & & \downarrow \\
 S^4 & \simeq & S(\mathbb{R} \oplus \mathbb{H})
 \end{array}
 \quad
 \begin{array}{c}
 (x, y) \\
 \downarrow \\
 \left(\begin{array}{l} t := 1 - |x|^2, \\ z := x \cdot y \end{array} \right)
 \end{array}
 \quad
 \begin{array}{l}
 |x|^2 + |y|^2 = 2 \\
 t^2 + |z|^2 = 1
 \end{array}$$

$$\pi_{\bullet}(S^4) \simeq \underbrace{\mathbb{Z}\langle \text{id}_{S^4} \rangle}_{\text{deg}=4} \oplus \underbrace{\mathbb{Z}\langle \text{Hopf}_{\mathbb{H}} \rangle}_{\text{deg}=7} \oplus \text{torsion}$$

Rational homotopy theory

disregards all torsion information:

$$\mathrm{Ho}(\mathrm{Spaces}_{\mathbb{Q}}) := \mathrm{Spaces} \left[\{ \text{Isos on rationalized homotopy groups} \}^{-1} \right]$$

$$\pi_k(S^n) \otimes \mathbb{Q} = \mathrm{Hom}_{\mathrm{Ho}(\mathrm{Spaces}_{\mathbb{Q}})} \left(S^k, S^n \right)$$

$$\pi_{\bullet}(S^4) \otimes \mathbb{Q} = \underbrace{\mathbb{Q}\langle \mathrm{id}_{S^4} \rangle}_{\mathrm{deg}=4} \oplus \underbrace{\mathbb{Q}\langle \mathrm{Hopf}_{\mathbb{H}} \rangle}_{\mathrm{deg}=7}$$

dg-Algebraic model for rational homotopy theory.

$\text{dgcAlg} := \{\text{differential graded-commutative algebras over } \mathbb{R}\}$

Homotopy theory of dg-algebras:

$$\begin{aligned} & \text{Ho}(\text{dgcAlg}^{\text{op}}) \\ & \quad := \\ & \text{dgcAlg}^{\text{op}} \left[\{\text{Isos on all cohomology groups}\}^{-1} \right] \end{aligned}$$

Quillen-Sullivan equivalence:

$$\begin{array}{ccc} \text{Ho}(\text{Spaces}) & \begin{array}{c} \xrightarrow{\mathcal{O}} \\ \perp \\ \xleftarrow{\mathcal{S}} \end{array} & \text{Ho}(\text{dgcAlg}^{\text{op}}) \\ \uparrow & & \uparrow \\ \text{Ho}(\text{Spaces}_{\mathbb{Q}, \text{nil}, \text{fin}}) & \begin{array}{c} \xrightarrow{\mathcal{O}} \\ \simeq \\ \xleftarrow{\mathcal{S}} \end{array} & \text{Ho}(\text{dgcAlg}_{\text{cn}, \text{fin}}^{\text{op}}) \end{array}$$

dg-Algebra model for the 4-sphere

Classical fact of rational homotopy theory:

$$\mathcal{O}(S^4) \simeq \mathbb{R}[\underbrace{\omega_4}_{\text{deg}=4}, \underbrace{\omega_7}_{\text{deg}=7}] / \left(\begin{array}{l} d\omega_4 = 0 \\ d\omega_7 = -\frac{1}{2}\omega_4 \wedge \omega_4 \end{array} \right)$$

This equation

$$dG_4 = 0$$

$$dG_7 + \frac{1}{2}G_4 \wedge G_4 = 0$$

is also the equations of motion in 11-dimensional supergravity for the M-brane flux.

Conjecture (Sati 13, Sect. 1.5):

The cohomology theory classifying M-branes is some flavour of degree-4 cohomotopy.

Super homotopy theory.

Via Quillen-Sullivan we may immediately generalize homotopy theory to *superspaces*:

$$\text{sdgcAlgebras} := \left\{ \begin{array}{l} \text{differential graded algebras} \\ \text{with} \\ \underbrace{\mathbb{Z}}_{\text{cohomological}} \times \underbrace{\mathbb{Z}/2}_{\text{super}} \text{-grading} \end{array} \right\}$$

In supergravity, cofibrant sdgc-algebras are known as “FDA”s.
We consider the corresponding homotopy theory:

$$\begin{aligned} & \text{Ho}(\text{SuperSpaces}_{\mathbb{Q}, \text{nil}, \text{fin}}) \\ & := \\ & \text{sdgcAlg}_{\text{cn}, \text{fin}}^{\text{op}} \left[\{ \text{Isos on all cohomology groups} \}^{-1} \right] \end{aligned}$$

see ncatlab.org/nlab/show/geometry+of+physics+-+superalgebra

Example: Super-Minkowski spacetimes.

The nilpotency condition on the fundamental group allows precisely the mild non-abelianness that goes with super-Minkowski spacetimes:

$$\mathcal{O} \left(\underbrace{\mathbb{T}^{p,1|\mathbb{N}}}_{\substack{\text{toroidal} \\ \text{super-Minkowski} \\ \text{spacetime}}} \right) := \mathbb{R} \left[\underbrace{\left(e^a \right)_{a=0}^p, \left(\psi^\alpha \right)_{\alpha=1}^N}_{\substack{\text{deg}=(1,\text{even}) \quad (1,\text{odd}) \\ \text{super vielbein}}} \right] / \underbrace{\left(\begin{array}{l} de^a = \bar{\psi} \Gamma^a \psi \\ d\psi^\alpha = 0 \end{array} \right)}_{\substack{\text{torsion} \\ \text{constraint}}}$$

Extension tower of super-Minkowski spacetime.

These appear in a sequence of iterated maximal R-symmetry invariant central extensions:

$$\begin{array}{c} \mathbb{R}^{10,1|\mathbf{32}} \\ \searrow \\ \mathbb{R}^{9,1|\mathbf{16}+\mathbf{16}} \longleftarrow \mathbb{R}^{9,1|\mathbf{16}} \Longrightarrow \mathbb{R}^{9,1|\mathbf{16}+\overline{\mathbf{16}}} \\ \searrow \\ \mathbb{R}^{5,1|\mathbf{8}} \Longrightarrow \mathbb{R}^{5,1|\mathbf{8}+\overline{\mathbf{8}}} \\ \swarrow \qquad \searrow \\ \mathbb{R}^{3,1|\mathbf{4}+\mathbf{4}} \longleftarrow \mathbb{R}^{3,1|\mathbf{4}} \\ \swarrow \qquad \searrow \\ \mathbb{R}^{2,1|\mathbf{2}+\mathbf{2}} \longleftarrow \mathbb{R}^{2,1|\mathbf{2}} \\ \swarrow \qquad \searrow \\ \mathbb{R}^0|\mathbf{1}+\mathbf{1} \longleftarrow \mathbb{R}^0|\mathbf{1} . \end{array}$$

The M2/M5-Brane cocycle.

On the top-dimensional 11d super Minkowski spacetime, there is a unique element in Spin(10, 1)-invariant rational cohomotopy in degree 4.

$$\mathbb{T}^{10,1|\mathbf{32}} \xrightarrow{\mu_{M2/M5}} S^4$$

$$\in \text{Ho}(\text{SuperSpaces}_{\mathbb{R}})^{\text{Spin}(10,1)}$$

In string/M theory this statement characterizes fundamental M2-brane and M5-branes via the WZW terms of their Green-Schwarz-type sigma-models.

2)
Parameterized stable homotopy theory
and
gauge enhancement

Central extensions are homotopy fibers.

That 11d super-spacetime is an extension
of 10d super-spacetime
means that it is a *homotopy fiber* of a 2-cocyle:

$$\begin{array}{ccc}
 \text{Ext}_{\overline{\psi}\Gamma_{10}\psi} \left(\mathbb{T}^{9,1} | \mathbf{16} + \overline{\mathbf{16}} \right) & \equiv & \mathbb{T}^{10,1} | \mathbf{32} \\
 & & \downarrow \text{homotopy fiber} \\
 & & \mathbb{T}^{9,1} | \mathbf{16} + \mathbf{16} \\
 & & \searrow \text{2-cocyle} \\
 & & \overline{\psi}\Gamma_{10}\psi \rightarrow BS^1
 \end{array}$$

The Ext/Cyc-Adjunction.

Proposition. The extension/homotopy fiber functor has a right adjoint

$$\mathrm{Ho}(\mathrm{Spaces}) \begin{array}{c} \xleftarrow{\mathrm{Ext}} \\ \perp \\ \xrightarrow{\mathrm{Cyc}} \end{array} \mathrm{Ho}\left(\mathrm{Spaces}/BS^1\right)$$

given by forming *cyclic loop spaces*:

$$\mathrm{Cyc}(X) := \mathrm{Maps}(S^1, X) // S^1$$

i.e. the homotopy quotient of the free loop space by the rigid rotation of loops.

Example: cyclification of the 4-Sphere

The cyclification of the 4-sphere is

$$\mathcal{O}(\text{Cyc}(S^4)) = \mathbb{Q}[h_3, h_7, \omega_2, \omega_4, \omega_6] / \left(\begin{array}{l} dh_7 = -\frac{1}{2}\omega_4 \wedge \omega_4 \\ \quad \quad \quad + \omega_2 \wedge \omega_6 \\ dh_3 = 0 \\ d\omega_{2p} = h_3 \wedge \omega_{2p-2} \end{array} \right)$$

Curiously, the terms in blue exhibit a truncation of rationalized twisted K-theory.

Below we will find also the rest of rationalized K-theory from the 4-sphere...

The Ext/Cyc-adjunct of $\mu_{M2/M5}$

Hence the M2/M5-cocycle

$$\mathbb{T}^{10,1|\mathbf{32}} \xrightarrow{\mu_{M2/M5}} S^4$$

induces its Ext/Cyc-adjunct

$$\begin{array}{ccc}
 \mathbb{T}^{9,1|\mathbf{16}+\overline{\mathbf{16}}} & \xrightarrow{\widetilde{\mu_{M2/M5}}} & \\
 \downarrow \eta & \searrow & \\
 \text{CycExt} \left(\mathbb{T}^{9,1|\mathbf{16}+\overline{\mathbf{16}}} \right) & \xrightarrow{\text{Cyc}(\mu_{M2/M5})} & \text{Cyc} (S^4) \\
 & \searrow & \swarrow \\
 & & B^2 S^1
 \end{array}$$

This is **double dimensional reduction** of M2/M5-cocycle
to the F1/NS5/D0/D2/D4 branes of type IIA string theory:

<p>Sullivan algebra of cyclified 4-sphere</p>	$dh_7 = -\frac{1}{2}\omega_4 \wedge \omega_4$ $+ \omega_2 \wedge \omega_6$ $dh_3 = 0$ $d\omega_{2p} = h_3 \wedge \omega_{2p-2}$
<p>Bianchi identities of NS1/NS5-flux and D($p \leq 4$) RR-fluxes:</p>	$dH_7 = -\frac{1}{2}F_4 \wedge F_4$ $+ F_2 \wedge F_6$ $dH_3 = 0$ $dF_{2p} = H_3 \wedge F_{2p-2}$

$$2p \in \{0, 2, 4\}$$

Extensions and actions

To make the double dimensional reduction more symmetric, we ask also S^4 to be an S^1 -extension

$$\begin{array}{ccc} \mathbb{R}^{10,1|\mathbf{32}} & & S^4 \\ \downarrow \text{homotopy fiber} & & \downarrow \text{homotopy fiber} \\ \mathbb{R}^{9,1|\mathbf{16}+\overline{\mathbf{16}}} & & S^4 // S^1 \\ & \searrow & \swarrow \\ & BS^1 & \end{array}$$

The base $S^4 // S^1$ of such an extension is necessarily the homotopy quotient by some S^1 -action on S^4 .

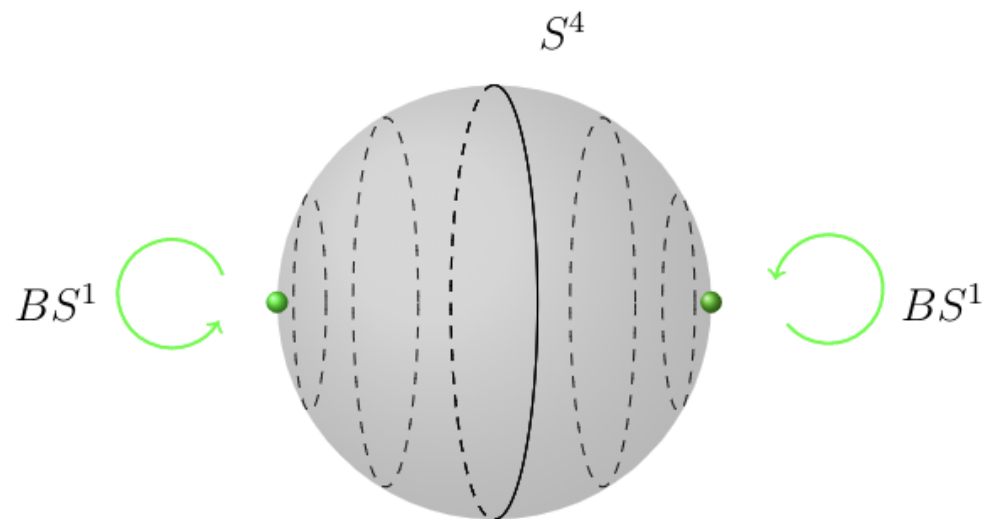
Example: Suspended Hopf action on S^4

The identifications

$$S^1 \simeq U(1) \subset SU(2) \simeq S(\mathbb{H})$$

$$S^4 \simeq S(\mathbb{R} \oplus \mathbb{H})$$

induce an S^1 -action on S^4 .



Attempted lift in Super homotopy theory

For any choice of S^1 -action on S^4
 we may ask for a lift of
 the double dimensional reduction of the M2/M5-cocycle:

$$\begin{array}{ccc}
 \mathbb{T}^{9,1|\mathbf{16}+\mathbf{16}} & \xrightarrow{\exists?} & S^4 // S^1 \\
 \downarrow \eta & \searrow \widetilde{\mu}_{M2/M5} & \downarrow \eta \\
 \text{CycExt} \left(\mathbb{T}^{9,1|\mathbf{16}+\overline{\mathbf{16}}} \right) & \xrightarrow{\text{Cyc} \left(\mu_{M2/M5} \right)} & \text{CycExt} \left(S^4 // S^1 \right) \\
 & \searrow & \swarrow \\
 & B^2 S^1 &
 \end{array}$$

Attempted lift in Super homotopy theory

For any choice of S^1 -action on S^4
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$$\begin{array}{ccc}
 \mathbb{T}^{9,1|\mathbf{16}+\mathbf{16}} & \xrightarrow{\exists} & S^4 // S^1 \\
 \downarrow \eta & \searrow \widetilde{\mu_{M2/M5}} & \downarrow \eta \\
 \text{CycExt} \left(\mathbb{T}^{9,1|\mathbf{16}+\overline{\mathbf{16}}} \right) & \xrightarrow{\text{Cyc} \left(\mu_{M2/M5} \right)} & \text{CycExt} \left(S^4 // S^1 \right) \\
 & \searrow & \swarrow \\
 & B^2 S^1 &
 \end{array}$$

Prop. *Such a lift does not exist, for any choice of action.
 But it may exist to first linear order.*

Stable homotopy theory

$$\mathrm{Ho}(\mathrm{Spectra}) := \mathrm{Spectra} \left[\{ \text{Isos on all stable homotopy groups} \}^{-1} \right]$$

Spectra stabilize the operation of forming loop spaces:

$$\begin{array}{ccc}
 \mathrm{Ho}(\mathrm{Spaces}^{*/}) & \begin{array}{c} \xleftarrow{\Sigma} \\ \xrightarrow{\Omega} \end{array} & \mathrm{Ho}(\mathrm{Spaces}^{*/}) \\
 \begin{array}{c} \Sigma^\infty \downarrow \\ \uparrow \Omega^\infty \end{array} & & \begin{array}{c} \Sigma^\infty \downarrow \\ \uparrow \Omega^\infty \end{array} \\
 \mathrm{Ho}(\mathrm{Spectra}) & \begin{array}{c} \xleftarrow{\Sigma} \\ \xrightarrow{\Omega} \end{array} & \mathrm{Ho}(\mathrm{Spectra}) \\
 & \text{with } \Sigma \approx \Omega &
 \end{array}$$

Loop spaces are the groups of homotopy theory.

Double loop space are the first-order commutative groups.

⋮

Hence: Spectra are the abelian groups of homotopy theory,

Hence: Σ^∞ is linearization in homotopy theory.

Parameterized stable homotopy theory

For $X \in \text{Ho}(\text{Spaces})$

$$\text{Ho}(\text{Spectra}_X) := \text{Spectra}_X \left[\left\{ \begin{array}{l} \text{Isos on all stable homotopy groups} \\ \text{for all homotopy fibers over } X \end{array} \right\}^{-1} \right]$$

Parameterized spectra stabilize forming homotopy-fiber wise loop spaces:

$$\begin{array}{ccc} \text{Ho} \left(\left(\text{Spaces}/X \right)^{X/} \right) & \begin{array}{c} \xleftarrow{\Sigma_X} \\ \xrightarrow{\Omega_X} \end{array} & \text{Ho} \left(\left(\text{Spaces}_X \right)^{X/} \right) \\ \begin{array}{c} \Sigma_X^\infty \downarrow \\ \uparrow \Omega_X^\infty \end{array} & & \begin{array}{c} \Sigma_X^\infty \downarrow \\ \uparrow \Omega_X^\infty \end{array} \\ \text{Ho} \left(\text{Spectra}/X \right) & \begin{array}{c} \xleftarrow{\Sigma_X} \\ \xrightarrow[\simeq]{\Omega_X} \end{array} & \text{Ho} \left(\text{Spectra}_X \right) \end{array}$$

Hence: Parameterized spectra are the bundles of abelian groups in homotopy theory.

Rational parameterized stable homotopy theory

Theorem (Braunack-Mayer 18) :

The Quillen-Sullivan dg-Model for rational homotopy theory generalizes to parameterized stable homotopy theory by modeling parameterized spectra by dg-modules:

$$\begin{array}{ccc}
 \mathrm{Ho} \left(\mathrm{Spectra}_{\mathcal{S}(A)} \right) & \begin{array}{c} \xrightarrow{\mathcal{O}} \\ \perp \\ \xleftarrow{\mathcal{S}} \end{array} & \mathrm{Ho} \left(\mathrm{dgMod}_A \right)^{\mathrm{op}} \\
 \uparrow & & \uparrow \\
 \mathrm{Ho} \left(\left(\mathrm{Spectra}_{\mathcal{S}(A)} \right)_{\mathbb{Q}, \mathrm{fin}} \right) & \begin{array}{c} \xrightarrow{\mathcal{O}} \\ \simeq \\ \xleftarrow{\mathcal{S}} \end{array} & \mathrm{Ho} \left(\mathrm{dgMod}_{A, \mathrm{bd}} \right)^{\mathrm{op}}
 \end{array}$$

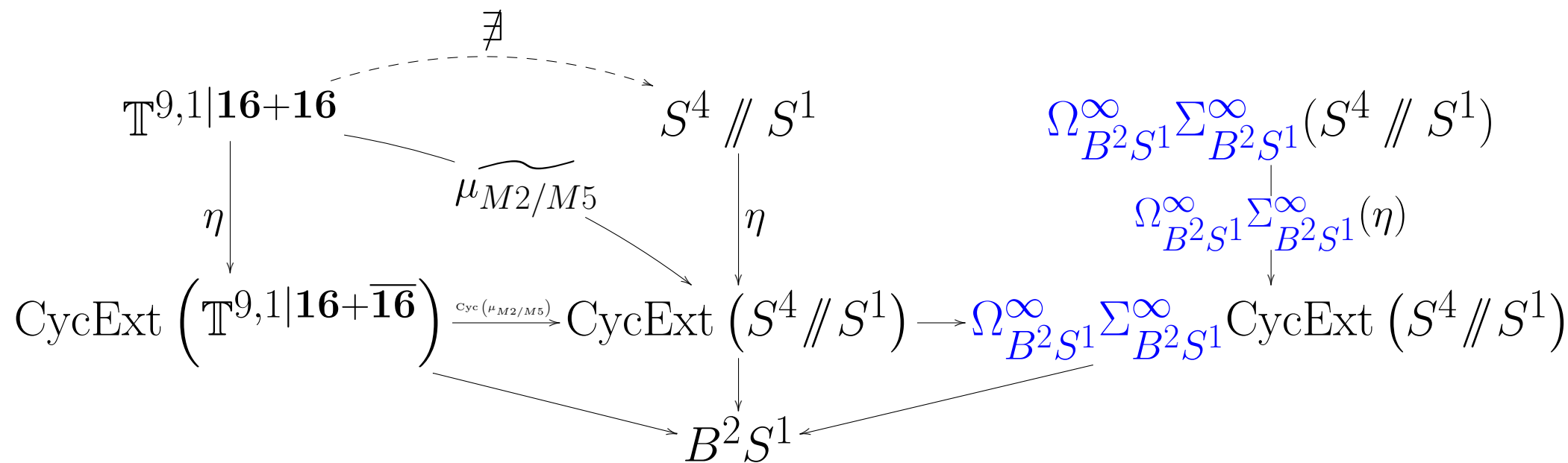
Linearized lift of Ext/Cyc-adjunct of M2/M5-Cocycle.

Remember that we were asking for a full lift,
which, however, does not exist:

$$\begin{array}{ccc}
 & \not\exists & \\
 & \curvearrowright & \\
 \mathbb{T}^{9,1|\mathbf{16}+\mathbf{16}} & & S^4 // S^1 \\
 \downarrow \eta & \xrightarrow{\widetilde{\mu}_{M2/M5}} & \downarrow \eta \\
 \text{CycExt} \left(\mathbb{T}^{9,1|\mathbf{16}+\overline{\mathbf{16}}} \right) & \xrightarrow{\text{Cyc}(\mu_{M2/M5})} & \text{CycExt} \left(S^4 // S^1 \right) \\
 & \searrow & \downarrow \\
 & & B^2 S^1
 \end{array}$$

Linearized lift of Ext/Cyc-adjunct of M2/M5-Cocycle.

But now we may **linearize** the coefficients:



Linearized lift of Ext/Cyc-adjunct of M2/M5-Cocycle.

Theorem: With the suspended Hopf S^1 -action on S^4 the linearized lifting problem has a unique solution:

$$\begin{array}{c}
 \mathbb{T}^{9,1|\mathbf{16}+\mathbf{16}} \xrightarrow{\mu_{F1/Dp}} \Omega_{B^2S^1}^\infty \Sigma_{B^2S^1}^\infty (S^4 // S^1) \\
 \downarrow \eta \quad \searrow \widetilde{\mu_{M2/M5}} \quad \downarrow \Omega_{B^2S^1}^\infty \Sigma_{B^2S^1}^\infty (\eta) \\
 \text{CycExt} \left(\mathbb{T}^{9,1|\mathbf{16}+\overline{\mathbf{16}}} \right) \xrightarrow{\text{Cyc}(\mu_{M2/M5})} \text{CycExt} (S^4 // S^1) \longrightarrow \Omega_{B^2S^1}^\infty \Sigma_{B^2S^1}^\infty \text{CycExt} (S^4 // S^1) \\
 \downarrow \quad \downarrow \quad \swarrow \\
 B^2S^1 \longleftarrow B^2S^1
 \end{array}$$

Linearized lift of Ext/Cyc-adjunct of M2/M5-Cocycle.

Theorem: This solution factors through a summand which is the (rationalized) **twisted K-theory spectrum**:

$$\begin{array}{ccccc}
 \mathbb{T}^{9,1|\mathbf{16}+\mathbf{16}} & \xrightarrow{\mu_{F1/Dp}} & \Omega_{B^2S^1}^{\infty-2}(\mathbf{ku} // BS^1) & \hookrightarrow & \Omega_{B^2S^1}^{\infty} \Sigma_{B^2S^1}^{\infty}(S^4 // S^1) \\
 \downarrow \eta & & \searrow \mu_{M2/M5} & & \downarrow \Omega_{B^2S^1}^{\infty} \Sigma_{B^2S^1}^{\infty}(\eta) \\
 \text{CycExt} \left(\mathbb{T}^{9,1|\mathbf{16}+\overline{\mathbf{16}}} \right) & \xrightarrow{\text{Cyc}(\mu_{M2/M5})} & \text{CycExt} (S^4 // S^1) & \longrightarrow & \Omega_{B^2S^1}^{\infty} \Sigma_{B^2S^1}^{\infty} \text{CycExt} (S^4 // S^1) \\
 & & \downarrow & & \swarrow \\
 & & B^2S^1 & &
 \end{array}$$

This is gauge enhancement of M2/M5-cocycle

to the full F1/Dp branes of type IIA string theory:

<p>dg-module for fiberwise stabilized suspended Hopf quotient of S^4 $\Sigma_{B^2 S^1}^\infty (S^4 // S^1)$</p>	<p>$dh_3 = 0$ $d\omega_{2p} = h_3 \wedge \omega_{2p-2}$</p>	<p>$2p$ \in \mathbb{Z}</p>
<p>Bianchi identities of NS1-flux and Dp RR-fluxes:</p>	<p>$dH_3 = 0$ $dF_{2p} = H_3 \wedge F_{2p-2}$</p>	

Beyond rational approximation?

We would like to eventually
solve the open problem
of lifting this discussion
beyond the rational approximation.

As a first step
towards solving this problem
we now lift the $S^1 \subset \mathrm{SU}(2)$ -action beyond rational
while keeping the spaces acted on rationally.

This hybrid approach
is
equivariant rational homotopy theory.

3)
Equivariant homotopy theory
and
black M-branes

Equivariant homotopy theory

For G a compact Lie group, the evident definition is:

$\text{Ho}(G\text{Spaces}) :=$

$G\text{CWComplexes} \left[\{G\text{-equivariant homotopy equivalences}\}^{-1} \right]$

The [equivariant Whitehead theorem](#)

relates this to fixed-point loci:

$\text{Ho}(G\text{Spaces}) \simeq$

$G\text{Spaces} \left[\left\{ \begin{array}{l} \text{Isos on all homotopy groups} \\ \text{after restriction to } H\text{-fixed points} \\ \text{for all closed subgroups } H \subset G \end{array} \right\}^{-1} \right]$

Systems of fixed point loci

But the H -fixed points are equivalently
the G -equivariant maps out of the orbit space G/H :

$$X^H = \text{Maps}(G/H, X)^G$$

Hence if we form the category of all possible G -orbit spaces

$$\text{Orb}_G := \{G/H\}_{H \subset G \text{ closed}} \subset G\text{Spaces}$$

Then the *system of fixed point loci* is extracted as

$$G\text{Spaces} \xrightarrow{Y} \text{PSh}(\text{Orb}_G, \text{Spaces})$$

$$X \mapsto \left(\begin{array}{ccc} G/H_1 & & X^{H_1} \\ \downarrow f & \mapsto & \uparrow \text{Maps}(f, X)^G \\ G/H_2 & & X^{H_2} \end{array} \right)$$

Equivariant homotopy theory is about fixed point loci.

Elmendorf's theorem: Equivariant homotopy theory is plain homotopy theory of the systems of fixed point loci:

$$\mathrm{Ho}(G\mathrm{Spaces}) \xrightarrow[\simeq]{Y} \mathrm{Ho}(\mathrm{PSh}(\mathrm{Orb}_G, \mathrm{Spaces}))$$

This induces in particular equivariance for non-classical homotopy theories:

Equivariant rational homotopy theory.

$$\mathrm{Ho}(G\mathrm{Spaces}_{\mathbb{Q}, \mathrm{nil}, \mathrm{fin}}) \simeq \mathrm{Ho}\left(\mathrm{PSh}\left(\mathrm{Orb}_G, \mathrm{dgcAlg}_{\mathrm{cn}, \mathrm{fin}}^{\mathrm{op}}\right)\right)$$

Equivariant rational **super homotopy theory.**

$$\mathrm{Ho}(G\mathrm{SuperSpaces}_{\mathbb{Q}, \mathrm{nil}, \mathrm{fin}}) \simeq \mathrm{Ho}\left(\mathrm{PSh}\left(\mathrm{Orb}_G, \mathrm{sdgcAlg}_{\mathrm{cn}, \mathrm{fin}}^{\mathrm{op}}\right)\right)$$

Example: $SU(2)$ -Action on $S^4_{\mathbb{Q}}$

Prop. The Hopf S^1 -action on S^4 is trivial on the rational homotopy type of S^4 , but in equivariant rational homotopy theory it is visible via its fixed point locus.

In fact this holds for the full $SU(2)$ -action:

$$\begin{array}{ccc}
 & SU(2)/1 & \\
 & \downarrow g & \\
 (S^4)^{(-)} : & SU(2)/1 & \longmapsto S^4 \\
 & \downarrow & \uparrow \text{id} \\
 & SU(2)/SU(2) & S^4 \\
 & & \uparrow \\
 & & S^0
 \end{array}$$

G_{HW} -action on S^4

There is also a \mathbb{Z}_2 -action on S^4 which does act non-trivially in rational homotopy theory, this is a reflection

$$\begin{array}{ccc}
 S^4 & \xrightarrow{\text{reverse}} & S^4 \\
 (x_1, \dots, x_4, x_5) & \longmapsto & (x_1, \dots, x_4, -x_5)
 \end{array}$$

The corresponding system of fixed points:

$$(S^4)^{(-)} \quad : \quad \begin{array}{ccc}
 \mathbb{Z}_2/1 & & S^4 \\
 \downarrow \sigma & & \uparrow \text{reverse} \\
 \mathbb{Z}_2/1 & \longmapsto & S^4 \\
 \downarrow & & \uparrow \\
 \mathbb{Z}_2/\mathbb{Z}_2 & & S^3
 \end{array}$$

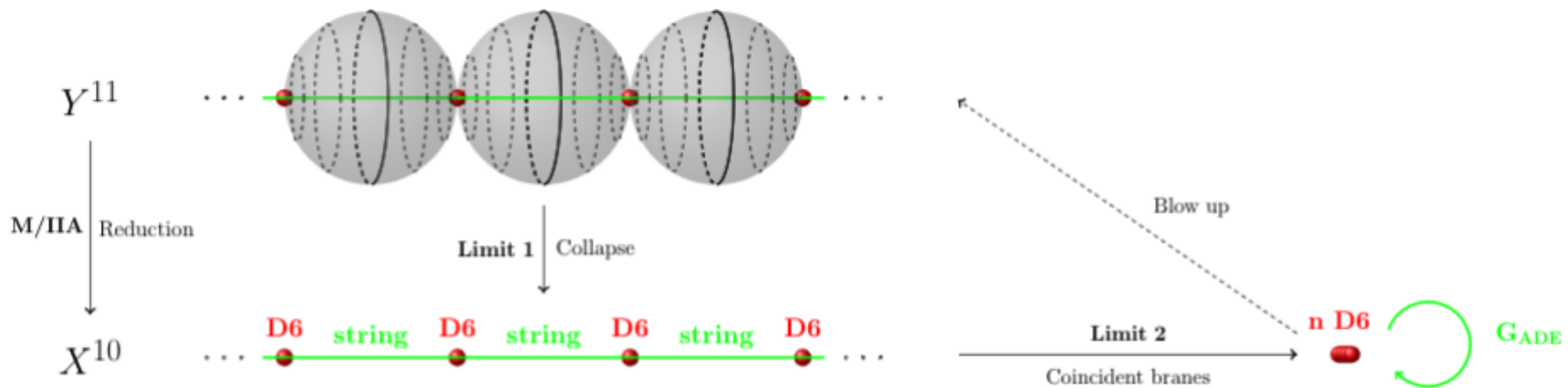
We will denote this action by G_{HW} .

ADE-classification of the finite subgroups of $SU(2)$:

Dynkin label	Finite subgroup $G_{\text{ADE}} \subset SU(2)$	Name of group
A_n	Z_{n+1}	Cyclic
D_{n+4}	$2D_{n+2}$	Binary dihedral
E_6	$2T$	Binary tetrahedral
E_7	$2O$	Binary octahedral
E_8	$2I$	Binary icosahedral

Resolution of ADE-singularities

Fact (du Val). The blow-up of an orbifold singularity fixed by a finite subgroup of $SU(2)$ is a system of spheres touching according to a Dynkin diagram:



Sen had suggested an interpretation in terms of M-brane physics. But the mathematical formulation of M-branes had remained an open problem.

Equivariant enhancement?

Hence we have lifted the 4-sphere to an object in $G_{\text{ADE}} \times G_{\text{HW}}$ -equivariant rational (super) homotopy theory.

Therefore we are now entitled to the following

Question: Does the M2/M5-cocycle have a corresponding equivariant enhancement?

$$\begin{array}{ccc} \mathbb{R}^{10,1|\mathbf{32}} \xrightarrow{\mu_{M2/M5}} S^4 & \in \text{Ho} (G_{\text{ADE,HW}} \text{SuperSpaces}_{\mathbb{R}}) & \\ & \downarrow \text{forget equivariance} & \\ \mathbb{R}^{10,1|\mathbf{32}} \xrightarrow{\mu_{M2/M5}} S^4 & \in \text{Ho} (\text{SuperSpaces}_{\mathbb{R}}) & \end{array}$$

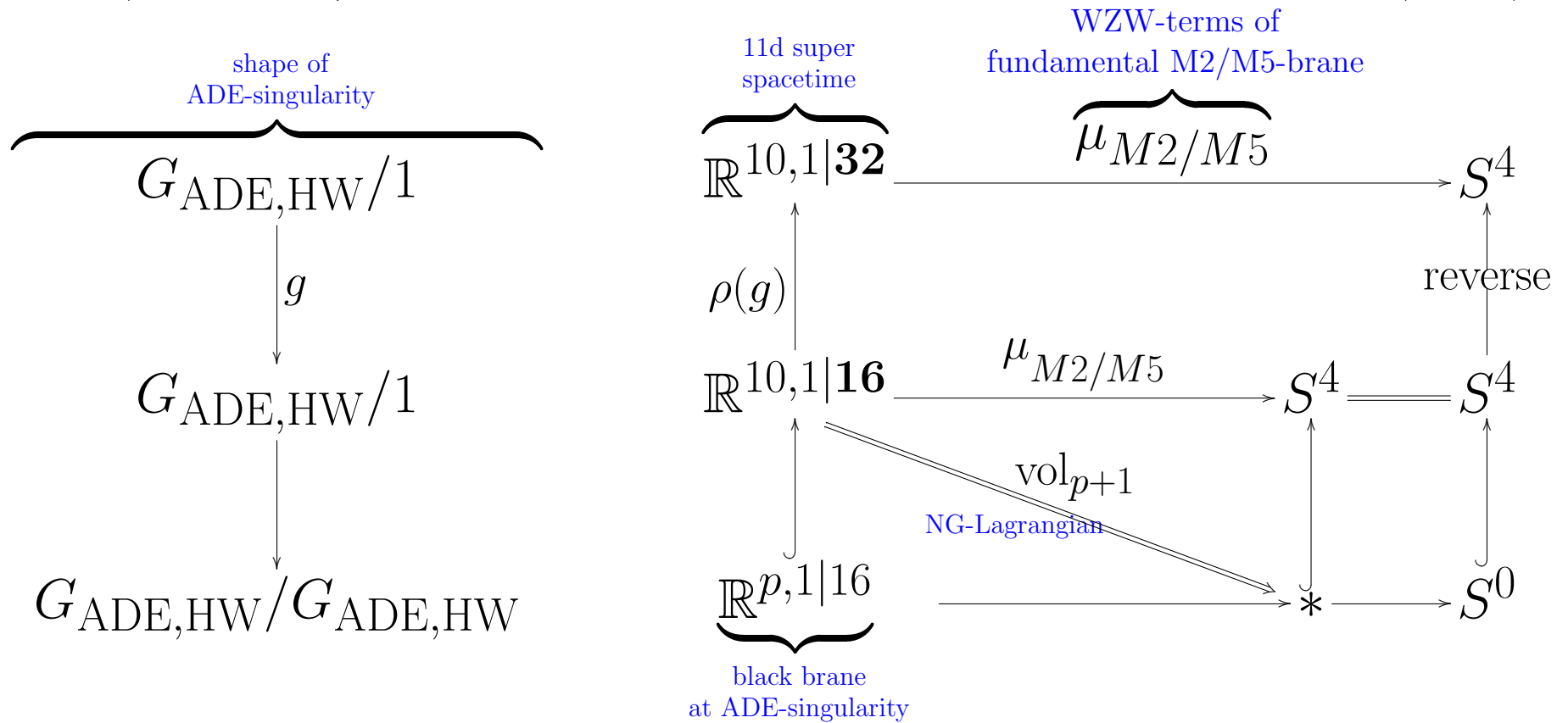
First, this requires identifying $G_{\text{ADE}} \times G_{\text{HW}}$ -actions on $\mathbb{R}^{10,1|\mathbf{32}} \dots$

Theorem: ADE-Singularities in 11d super spacetime.

Group action on $\mathbb{R}^{10,1 32}$	Possible singular locus					BPS
type H	NS1 _H	M2	$1/2$ M5 _H	MKK6	M9 _H	
type I	E1		$1/2$ M5 _I		M9 _I	
G_{ADE}				$\mathbb{R}^{6,1 16}$		$1/2$
G_{ADE}		$\mathbb{R}^{2,1 \geq 8 \cdot 2}$				$\geq 1/2$
$G_{\text{HW}} = \mathbb{Z}_2$					$\mathbb{R}^{9,1 16}$	$1/2$
$G_{\text{ADE,HW}} = \mathbb{Z}_2$	$\mathbb{R}^{1,1 16 \cdot 1}$		$\mathbb{R}^{5,1 8+8}$			$1/2$
$G_{\text{ADE}} \times G_{\text{HW}}$	$\mathbb{R}^{1,1 8 \cdot 1}$		$\mathbb{R}^{5,1 8}$			$1/4$

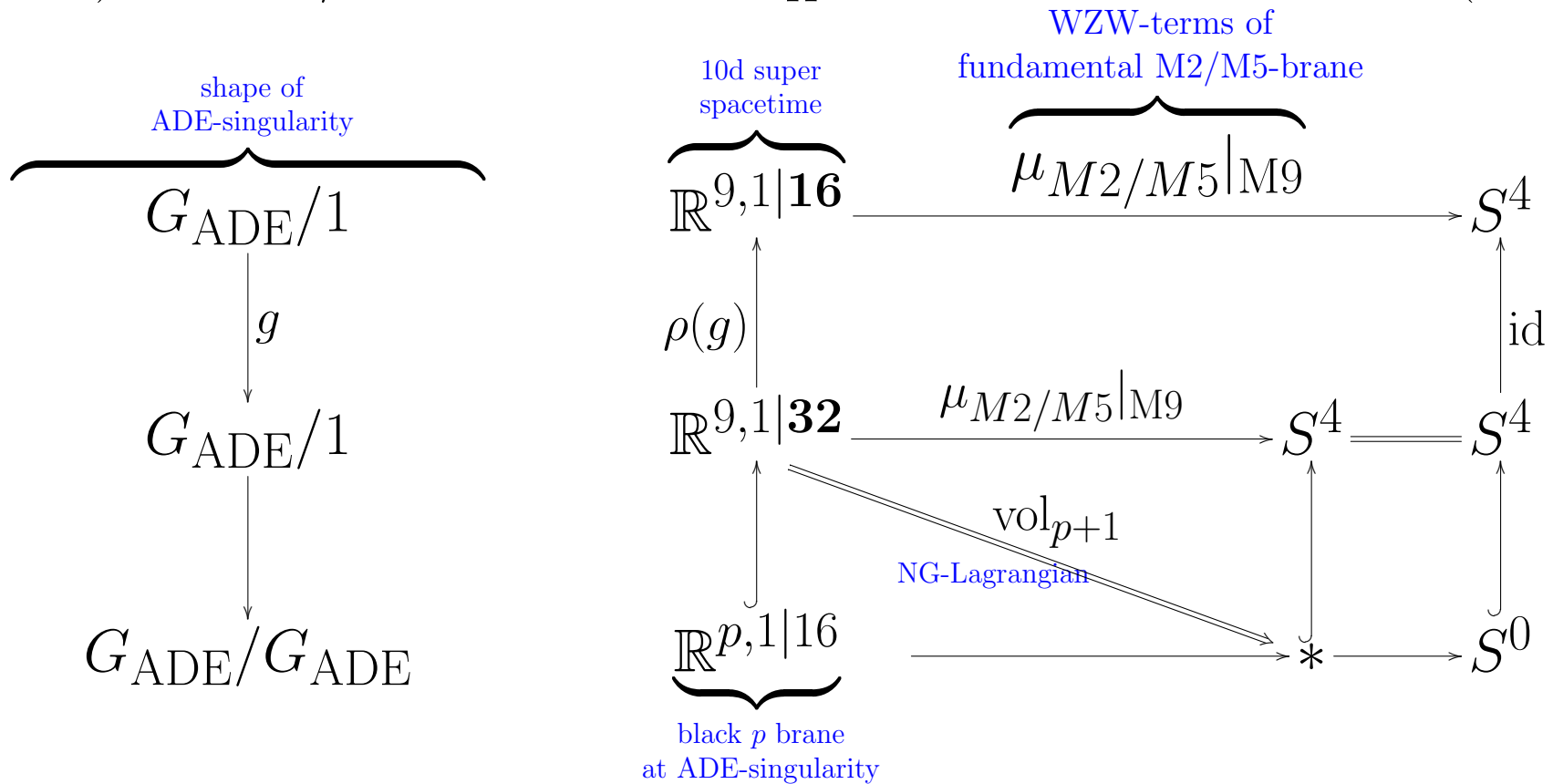
Theorem: There are two $G_{\text{ADE,HW}}$ -equivariant enhancements of the M2/M5-brane cocycle, exhibiting...

- 1) ... the 1/2-BPS black $1/2\mathbf{M5}$ of $D = 5 + 1$ and $N = (2, 0)$
- 2) ... the 1/2-BPS black $\mathbf{NS1}_H$ of $D = 1 + 1$ and $N = (16, 0)$



Theorem: There are two G_{ADE} -equivariant enhancements of the M2/M5-brane cocycle restricted to $M9_H$, exhibiting...

- 1) ... the $1/4$ -BPS black $1/2\mathbf{M5}$ of $D = 5 + 1$ and $N = (1, 0)$
- 2) ... the $1/4$ -BPS black $\mathbf{NS1}_H$ of $D = 1 + 1$ and $N = (8, 0)$



Conclusion:

A fair bit of
the expected structure of **M-theory**
emerges out of the superpoint $\mathbb{R}^{0|1}$
in rational equivariant super homotopy theory.

Evident Conjecture:

The full theory emerges
once passing beyond the rational approximation
in **full super-geometric homotopy** theory.
(arXiv:1310.7930).

Epilogue

In full super-geometric homotopy theory
the superpoint $\mathbb{R}^{0|1}$ itself
emerges from \emptyset

$$\begin{array}{ccccccc}
 \text{id} & \dashv & \text{id} & & & & \\
 \vee & & \vee & & & & \\
 \rightrightarrows & \dashv & \rightsquigarrow & \dashv & \boxed{\mathbb{R}^{0|1}} & & \\
 & & \vee & & \vee & & \\
 & & \mathfrak{R} & \dashv & \boxed{\mathbb{D}} & \dashv & \text{Et} \\
 & & & & \vee & & \vee \\
 & & & & \boxed{\mathbb{R}} & \dashv & \mathfrak{b} & \dashv & \# \\
 & & & & & & \vee & & \vee \\
 & & & & & & \emptyset & \dashv & *
 \end{array}$$

(Schreiber 16, FOMUS proceedings)