

16. By applying the definitions, we see that, for all $f \in C^\infty(N)$, we have

$$\phi_*(\gamma'(t))(f) = \gamma'(t)(f \circ \phi) = \frac{d}{dt}f(\phi(\gamma(t))) = \frac{d}{dt}f(\phi \circ \gamma(t)) = (\phi \circ \gamma)'(t)(f)$$

17. For $f, g \in C^\infty(N)$ and $\alpha \in \mathbb{R}$, we have, by linearity of $v \in T_pM$, that

$$(\phi_*v)(f + \alpha g) = v((f + \alpha g) \circ \phi) = v(f \circ \phi) + \alpha v(g \circ \phi) = (\phi_*v)f + \alpha(\phi_*v)(g)$$

18. For $v \in \text{Vect}(M)$ and $\phi : M \rightarrow N$ a diffeomorphism, we may define $\phi_*v \in \text{Vect}(N)$ as follows: For $q \in N$ and $f \in C^\infty(N)$, define

$$(\phi_*v)(f)(q) = v(f \circ \phi)(\phi^{-1}(q)) \quad \text{i.e.} \quad (\phi_*v)_{\phi(p)}(f) := v_p(\phi^*f)$$

(where $f \in C^\infty(N)$ and $v_p(g) := v(g)(p)$, etc.) To show that $\phi_*v \in \text{Vect}(N)$, note that

$$(\phi_*v)_{\phi(p)}(fg) = v_p((fg) \circ \phi) = f(\phi(p))(\phi_*v)(g) + (\phi_*v)(f)g(\phi(p))$$

to obtain the Leibniz property. Linearity follows similarly.

19. Note that $\phi(x, y) = (x \cos \theta - y \sin \theta, x \sin \theta + y \cos \theta)$, and so

$$\begin{aligned} ((\phi_*\partial_x)f, (\phi_*\partial_y)f) &= (\partial_x(f \circ \phi), \partial_y(f \circ \phi))|_{\phi(x,y)} \\ &= D(f \circ \phi)|_{\phi(x,y)} \\ &= Df|_{\phi(x,y)}D\phi|_{(x,y)} \\ &= (\partial_x f, \partial_y f)|_{\phi(x,y)} \begin{pmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{pmatrix} \\ &= (\cos \theta \partial_x f + \sin \theta \partial_y f, -\sin \theta \partial_x f + \cos \theta \partial_y f) \end{aligned}$$

where D denotes the derivative as linear operator, given by the Jacobian matrix.

20. If $\gamma : \mathbb{R} \rightarrow \mathbb{R}^2 : t \mapsto (x(t), y(t))$ is an integral curve, we want that $\gamma'(t) = x(t)^2 \partial_x + y(t) \partial_y$, i.e. that

$$\partial_x f \frac{dx}{dt} + \partial_y f \frac{dy}{dt} = \frac{d}{dt}f(x(t), y(t)) = \gamma'(t)(f) = x(t)^2 \partial_x f + y(t) \partial_y f$$

for all $f \in C^\infty(\mathbb{R}^2)$. So e

$$\frac{dx}{dt} = x(t)^2 \quad \frac{dy}{dt} = y(t)$$

and hence either $x(t) = 0$ or

$$x(t) = \frac{1}{A-t} \quad y(t) = Be^t$$

for constants A, B . These blow up at $t = A$.

21. The semigroup property is intuitively clear from the “water” analogy in BM: If a water molecule is at point p at time zero, it will be at $\phi_t(p)$ at time t . s units of time later, i.e. at time $t + s$ the point particle will be at $\phi_s(\phi_t(p))$ — this assumes that the vector field is constant over time. Doing this in one step, the particle will also be at $\phi_{t+s}(p)$ at time $t + s$.

Actually, though intuitively clear, I didn’t find the semigroup property all that easy to prove (semi)–formally, partly because BM aren’t very precise, and partly because there’s probably an easier way I am not seeing. Anyway,

- (i) If v is a vector field on a manifold M , there is at a unique integral curve $\gamma_p(t)$ through each point $p \in M$ such that $\gamma_p(0) = p$. This is implied by a statements in BM (“Let $\phi_t(p)$ be **the** integral curve of v through $p \in M$ ”), and “Show that ϕ_0 is the identity map”, but presumably needs some Lipschitz-like conditions, as for ODE’s.
- (ii) Certainly the integral curve as function is not unique, but nor is it sufficient to simply specify the range of an integral curve as a subset of M . The speed at which the curve is traversed is important, but the starting point is not. To be more precise, note that if $\gamma(t)$ is an integral curve through a point p , and $s \in \mathbb{R}$ is fixed, then $\bar{\gamma}(t) := \gamma(t+s)$ is also an integral curve through p (but starting at a different point). Indeed, if $f \in C^\infty(M)$, and $h_s(t) := t+s$, then by the chain rule

$$(f \circ \bar{\gamma})'(t) = (f \circ \gamma \circ h_s)'(t) = (f \circ \gamma)'(h_s(t))h_s'(t) = (f \circ \gamma)'(t+s)$$

Hence $\bar{\gamma}'(t) = \gamma'(t)$ as tangent vectors in $T_{\gamma(t+s)}M$, and hence $\bar{\gamma}$ is also an integral curve through p . Yet $\bar{\gamma} \neq \gamma$ as curves.

- (iii) For $p \in M$, let $\psi_p(t) := \phi_t(p)$, where $\{\phi_t : M \rightarrow M\}$ is the flow generated by a given vector field v , with $\phi_0 = \text{id}_M$. Then by the previous bullet point, for fixed s , we have that $\psi_p(t+s)$ and $\psi_p(t)$ are the “same” integral curve, namely the integral curve through p . However $\psi_p(t+s)$ starts at the point $\psi_p(s)$ at 0, as does the integral curve $\psi_{\psi_p(s)}(t)$. By uniqueness, it follows that $\psi_p(t+s) = \psi_{\psi_p(s)}(t)$. Going back from ψ ’s to ϕ ’s, we see that

$$\phi_{t+s}(p) = \psi_p(t+s) = \psi_{\psi_p(s)}(t) = \phi_t(\phi_s(p))$$

i.e. $\phi_{t+s} = \phi_t \circ \phi_s$.

22. $v = \partial_r$, $w = \frac{1}{r}\partial_\theta$, so

$$[v, w] = \partial_r(\frac{1}{r}\partial_\theta) - \frac{1}{r}\partial_\theta\partial_r = -\frac{1}{r^2}\partial_\theta = -\frac{w}{r} = -\frac{y\partial_x - x\partial_y}{x^2 + y^2}$$

- 23.

$$vw(f)(p) = \frac{d}{dt}w(f)(\phi_t(p))\Big|_{t=0} = \frac{d}{dt}\frac{d}{ds}f(\psi_s(\phi_t(p)))\Big|_{t,s=0}$$

24. 1), 2) are easy.

The Jacobi identity can be proved as follows: Multiplying out the lefthand side $[u, [v, w]] + [v, [w, u]] + [w, [u, v]]$ we will obtain 12 terms. Each such term can be regarded as a permutation of uvw — of which there are $3!=6$ in total — and each term will occur exactly twice, but with opposite signs. Thus $+abc$ will occur in the expansion of $[a, [b, c]]$, and $-abc$ in that of $[c, [a, b]]$. But this “explanation” is more complicated than simply multiplying the whole thing out:

$$\begin{aligned} [u, [v, w]] &= uvw - uww - vwu + wvu \\ [v, [w, u]] &= vwu - vuv - wuv + uvw \\ [w, [u, v]] &= wuv - wvu - uvw + vuv \end{aligned}$$

Adding, we see that all terms cancel.

25. For $\omega, \mu \in \Omega^1(M)$, $v, w \in \text{Vect}(M)$ and $f, g \in C^\infty(M)$, we have

$$(\omega + \mu)(v + w) = \omega(v + w) + \mu(v + w) = \omega(v) + \omega(w) + \mu(v) + \mu(w) = (\omega + \mu)(v) + (\omega + \mu)(w)$$

and

$$(\omega + \mu)(gv) = \omega(gv) + \mu(gv) = g\omega(v) + g\mu(v) = g(\omega + \mu)(v)$$

Also

$$(f\omega)(v + w) = f(\omega(v + w)) = f(\omega(v) + \omega(w)) = (f\omega)(v) + (f\omega)(w)$$

and

$$(f\omega)(gv) = f(\omega(gv)) = f(g\omega(v)) = g(f\omega)(v)$$

26. We have

$$\begin{aligned} f(\omega + \mu)(v) &= f(\omega(v) + \mu(v)) = f\omega(v) + f\mu(v) \\ (f + g)(\omega)(v) &= f\omega(v) + g\omega(v) \\ (fg)(\omega)(v) &= f(g\omega(v)) \\ 1\omega(v) &= \omega(v) \end{aligned}$$

27.

$$\begin{aligned} d(f + g)(v) &= v(f + g) = v(f) + v(g) = df(v) + dg(v) \\ d(\alpha f)(v) &= v(\alpha f) = \alpha v(f) = \alpha df(v) \\ (f + g)dh(v) &= (f + g)v(h) = fv(h) + gv(h) = f dh(v) + g dh(v) \\ d(fg)(v) &= v(fg) = gv(f) + fv(g) = g df(v) + f dg(v) \end{aligned}$$

28. If $v \in \text{Vect}(\mathbb{R}^n)$, then $v = v^i(x)\partial_i$ for some $v^i \in C^\infty(\mathbb{R}^n)$. Note that we have $dx^j(v) = v^i(x)\partial_i x^j = v^j(x)$ so that

$$df(v) = v(f) = v^i(x)\partial_i f = \partial_i f dx^i(v)$$

29. Suppose that $\omega := \omega_\mu dx^\mu = 0$, and define $v := \omega^\nu \partial_\nu$, where $\omega^\nu := \omega_\nu$ (i.e. equal as functions in $C^\infty(\mathbb{R}^n)$). Then

$$0 = \omega(v) = \omega_\mu dx^\mu(\omega^\nu \partial_\nu) = \omega_\mu \omega_\nu \delta_\nu^\mu = \omega_\mu \omega^\mu = \sum_{\mu=1}^n (\omega_\mu)^2$$

30. I can't do this (yet) without coordinates, i.e. without assuming every 1-form is locally of the form $\omega = \omega_\mu dx^\mu$. Assuming that, it is easy.

31. $(\text{id}^*\omega)(v) = \omega(\text{id}(v)) = \omega(v)$.

$$(gf)^*(\omega)(v) = \omega(g(f(v))) = (g^*\omega)(f(v)) = (f^*g^*\omega)(v).$$

32. For $\phi : M \rightarrow N$, we put $\phi^* : \Omega^1(N) \rightarrow \Omega^1(M)$ by $\phi^*(\omega)(v)(p) = \omega(\phi_*(v))(\phi(p))$. Let $f \in C^\infty(M)$, $v \in \text{Vect}(M)$, and put $q := \phi(p)$. Note that $\phi_*(fv)_p(g) = (fv)_p(g \circ \phi) = f(p)v_p(g \circ \phi) = (f\phi_*(v))_p(g)$, so that $\phi_*(fv) = f\phi_*(v)$. Hence

$$\phi^*(\omega)_p(fv)_p = \omega_q(\phi_*(fv)_p) = \omega_q(f\phi_*(v)_p) = f(p)\omega_q(\phi_*(v)_p) = (f\phi^*(\omega))_p(v_p)$$

so that indeed $\phi^*(\omega)(fv) = f\phi^*(\omega)(v)$.

33. Note first that $dx = d\text{id}_{\mathbb{R}}$. Hence

$$\phi^*(dx)_t(v) = dx_{\text{id}(t)}(\phi_*(v)) = \phi_*(v)_t(\text{id}) = v_t(\text{id} \circ \phi) = v_t(\phi) = d\phi_t(v)$$

i.e. $\phi^*(dx) = d\phi$. But clearly $d\phi_t = \frac{\partial \phi}{\partial t} dt = \cos t dt$.

34. $\phi(x, y) := \begin{pmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} = (x \cos \theta - y \sin \theta, x \sin \theta + y \cos \theta)$, so with π_1 denoting the projection onto the first component, we see that

$$\phi^*(dx) = \phi^*(d\pi_1) = d(\phi^*\pi_1) = d(\pi_1 \circ \phi) = d(x \cos \theta - y \sin \theta) = \cos \theta dx - \sin \theta dy$$

using linearity of d .

35. Let $\varphi : U \rightarrow \mathbb{R}^n$ be a chart. We are writing x^μ instead of $x^\mu \circ \varphi = \varphi^*(x^\mu)$, ∂_μ instead of $(\varphi^{-1})_*(\partial_\mu)$, and dx^μ instead of $\varphi^*(dx^\mu)$. We have to check that

$$d(\phi^* x^\mu) = \varphi^*(dx^\mu)$$

but we already know that.

36. By definition, $dx^\lambda(\partial_\mu) = \partial_\mu(x^\lambda) = \delta_\mu^\lambda$. Thus if $dx'^\nu = S_\mu^\nu dx^\mu$, we see that $dx'^\nu(\partial_\lambda) = S_\lambda^\nu$. But $dx'^\nu(\partial_\lambda) = \partial_\lambda x'^\nu$ and hence $S_\mu^\nu = \frac{\partial x'^\nu}{\partial x^\mu}$. Then if $\omega'_\nu dx'^\nu = \omega_\mu dx^\mu$, we see that

$$\omega'_\nu \frac{\partial x'^\nu}{\partial x^\mu} dx^\mu = \omega_\mu dx^\mu$$

from which $\omega_\mu = \frac{\partial x'^\nu}{\partial x^\mu} \omega'_\nu$. Now to obtain $\omega'_\nu = \frac{\partial x^\mu}{\partial x'^\nu} \omega_\mu$ either interchange primed and non-primed coordinates and μ and ν , or else multiply both sides by $\frac{\partial x^\lambda}{\partial x'^\nu}$ (with summation convention), noting that $\frac{\partial x^\lambda}{\partial x'^\nu} \frac{\partial x'^\nu}{\partial x^\mu} = \frac{\partial x^\lambda}{\partial x^\mu} = \delta_\mu^\lambda$.

37. On the left we have

$$\phi^*(dx'^\nu)(\partial_\lambda) = d(\phi^* x'^\nu)(\partial_\lambda) = \partial_\lambda(\phi^* x'^\nu) = \text{sloppy } \frac{\partial x'^\nu}{\partial x^\lambda}$$

and on the right we obtain the same:

$$\frac{\partial x'^\nu}{\partial x^\mu} dx^\mu(\partial_\lambda) = \frac{\partial x'^\nu}{\partial x^\mu} \delta_\lambda^\mu = \frac{\partial x'^\nu}{\partial x^\lambda}$$